



On the two heating rate method of determination of the activation energy of TL peaks under different heating schemes

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Abstract

In the present paper, we consider the application of the Booth's method and another variant of two heating rate method suggested by Pagonis. The Booth's method is based on the shift of the TL peak temperature with change of heating rate and another variant is related with shift of peak intensity with heating rate. We have considered the applicability of the method for non- first order kinetic peaks with temperature dependent frequency factor.

Key words: Thermoluminescence, activation energy, two-heating rates method, proportional error

1. INTRODUCTION

Thermoluminescence (TL) is a well-established technique that is widely applied in dosimetry, archeological dating and for detecting changes in defect concentration in insulators [1,2]. The underlying premise of TL emission is that the exposure of a material to the ionizing radiation causes a redistribution of charges in the defect centres within the material. When the material is heated at a controlled heating scheme, the TL is emitted as a temperature dependent set of peaks collectively known as a glow curve. Activation energy (E) is one of the important parameters connected with a typical peak of a glow curve. The stability of a peak is decided by its activation energy which is crucial for dating purposes.

Booth method [3] is a particular type of the simple two different heating rates method [1,2] for the determination of activation energy (E) of TL peak. Recently Pagonis et al [4] suggested another variant of two-heating rates method hereafter called Pagonis method for the determination of activation energy for TL peaks. In the present paper we plan to

adjudge the suitability of the Booth and Pagonis method for the determination of activation energy of a TL peak.

2. METHODOLOGY

The TL intensity $I(T)$ of a first order ($b=1$) TL peak is given by

$$I(T) = sn_0 \exp\left(-\frac{E}{kT}\right) \exp\left[-\frac{s}{\beta} \int_{T_0}^T \exp\left(-\frac{E}{kT'}\right) dT'\right] \quad (1)$$

whereas for a non-first order TL peak one can write

$$I(T) = sn_0 \exp\left(-\frac{E}{kT}\right) \times \left[1 + (b-1) \frac{s}{\beta} \int_{T_0}^T \exp\left(-\frac{E}{kT'}\right) dT'\right]^{-\frac{b}{b-1}} \quad (2)$$

where n_0 is the initial number of trapped electrons, s is the of frequency factor and b is the order of kinetics. β is the heating rate corresponding to the linear heating scheme

$$T = T_0 + \beta t \quad (3)$$

where T is the temperature at any time t and T_0 is the initial temperature at $t=0$, k is the Boltzmann constant

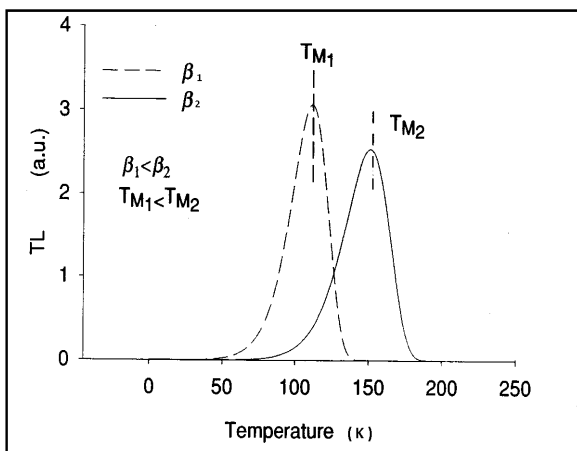


Fig. 1 : Schematic diagram of two-heating rate method

The peak temperature T_m corresponding to maximum intensity I_m of a TL peak for $b=1$ (Eqn. 1) is given by

$$\frac{\beta E}{kT_m^2} = s \exp\left(-\frac{E}{kT_m}\right) \quad (4)$$

Whereas for general order kinetics the corresponding maximum intensity condition can be expressed as

$$\left[1 + (b-1) \frac{s}{\beta} \int_{T_0}^{T_m} \exp\left(-\frac{E}{kT'}\right) dT'\right] = \frac{bskT_m^2}{\beta E} \exp\left(-\frac{E}{kT_m}\right) \quad (5)$$

Now using the asymptotic expression of the integral

$$\int_{T_0}^{T_m} \exp\left(-\frac{E}{kT}\right) dT \quad \text{one can write [4]}$$

$$\int_{T_0}^{T_m} \exp\left(-\frac{E}{kT}\right) dT = \frac{kT_m^2}{E} \exp\left(-\frac{E}{kT_m}\right) \left[1 - \frac{2kT_m}{E}\right] \quad (6)$$

Whereas without any loss of generality T_0 may be replaced by 0 [5] so that eqn. (4) can be expressed as

$$\frac{\beta E}{kT_m^2} = s \exp\left(-\frac{E}{kT_m}\right) \left[1 + (b-1) \frac{2kT_m}{E}\right] \quad (7)$$

It can be easily shown that for experimental TL peaks, the second term within the parenthesis can be neglected compared to unity so that both for first order and non-first order peaks eqn.(4) for maximum can be expressed as

$$\frac{\beta E}{kT_m^2} = s \exp\left(-\frac{E}{kT_m}\right) \quad (8)$$

If we consider two different heating rates (β_1, β_2) activation energy as calculated by Booth method is given by [3], where T_{m1} and T_{m2} are corresponding peak temperatures (Fig. 1).

$$E_B = k \frac{T_{m1} T_{m2}}{T_{m1} - T_{m2}} \ln \left[\left(\frac{\beta_1}{\beta_2} \right) \left(\frac{T_{m2}}{T_{m1}} \right)^2 \right] \quad (9)$$

Using equation (1) and (2) the expression for $\ln I_m$ for $b=1$ and $b \neq 1$ are given by [4]

$$\ln I_m = -\frac{E}{kT_m} + \ln sn_0 - \frac{s}{\beta} \frac{kT_m^2}{E} \left[1 - \frac{2kT_m}{E}\right] \quad b=1 \quad (10)$$

$$\ln I_m = -\frac{E}{kT_m} + \ln sn_0 - \frac{b}{b-1} \ln \left[1 + (b-1) \frac{s}{\beta} \frac{kT_m^2}{E} \left(1 - \frac{2kT_m}{E} \right) \right] \quad b \neq 1 \quad (11)$$

It has been shown by Pagonis et al [4] that the third term in equations (10) and (11) can be neglected in comparison with the first two terms so that both equations (10) and (11) can be expressed as

$$\ln I_m \cong -\frac{E}{kT_m} + \ln sn_0 \quad (12)$$

Now we consider two different heating rates (β_1, β_2) and I_{m1} and I_{m2} are the corresponding peak intensities. The activation energy as calculated by the method proposed by Pagonis et al [4]

$$E_p = \frac{kT_{m1}T_{m2}}{T_{m1} - T_{m2}} \ln \frac{I_{m1}}{I_{m2}} \quad (13)$$

In case of temperature dependent frequency factor, where

$$s = s_0 T^a \quad \text{with } -2 \leq a \leq 2 \quad (14)$$

the TL intensity of a glow peak is given by

$$I(T) = n_0 s_0 T^a \exp(-E/kT) \times \exp \left(-\frac{E}{kT} - (s_0/\beta) \int_{T_0}^T T'^a \exp(-E/kT') dT' \right) \quad b=1 \quad (14)$$

$$I(T) = n_0 s_0 \exp(-E/kT) \times \left\{ 1 + \frac{(b-1)s_0}{\beta} \int_{T_0}^T T'^a \exp(-E/kT') dT' \right\}^{-\frac{b}{b-1}} \quad b \neq 1 \quad (15)$$

The maxima condition for TL peaks corresponding to the first order and non-first order kinetics are given respectively by

$$\frac{\beta E}{kT_m^2} - \frac{s}{\beta} T_m^a \exp\left(-\frac{E}{kT_m}\right) + \frac{a}{T_m} = 0 \quad (16)$$

And

$$\frac{\beta E}{kT_m^2} + \frac{a}{T_m} = \frac{bs_0 T_m^a \exp(-E/kT_m)}{\left[1 + s_0 \frac{(b-1)}{\beta} \int_{T_0}^{T_m} T'^a \exp(-E/kT') dT' \right]} \quad (17)$$

Equations (16) and (17) can be used to determine peak temperature T_m and eqn. (12) is used to determine maximum intensity I_m .

3. RESULTS AND DISCUSSION

In order to compute the errors involved in different methods such as proposed by Booth [3] and Pagonis et al [4] we have computed T_m and I_m using relevant equations (14-17) for particular values of E , s and β . Using the computed values of T_m and I_m for two heating rates we have computed activation energies E_B and E_p by using respectively the method of Booth [3] and Pagonis et al [4]. After computing E_B and E_p , we compute the percentile

proportional error $\delta_B (= \frac{|E - E_B|}{E} \times 100\%)$ and $\delta_P (= \frac{|E - E_P|}{E} \times 100\%)$ in determination of activation energy by Booth[3] and Pagonis method[4]. The plot of δ_B and δ_P with variation of b ($1 \leq b \leq 2.5$) for $a=-2,0,2$ and for $E=1$ eV, $s_0 = 10^{12}$ units is depicted in Figure 2.

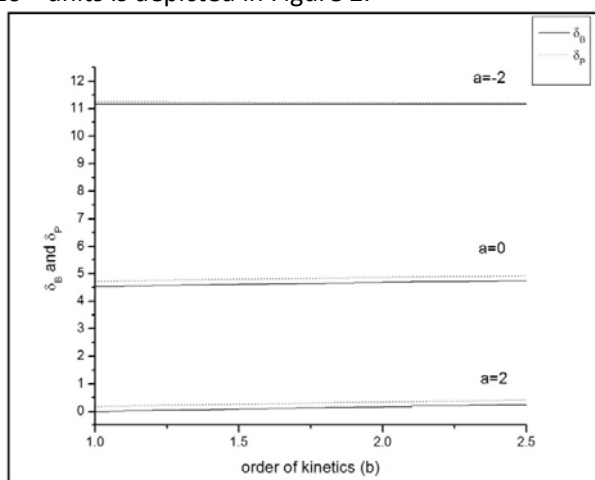


Fig 2. Variation of δ_B and δ_P with order of kinetics for different values of ‘a’ in Computer simulated glow curve (E=1 eV, s=10¹² units)

In all the cases δ_B and δ_P is constant with b . δ_P is marginally larger than δ_B . We have run computer codes for a value of E ranging from 0.1 eV to 1.6 eV. δ_B and δ_P in all the cases have been found insensitive to the input value of E used in the calculation. In Figure 3, we present the variation of δ_B and δ_P as a function of $\log s_0$ for $b = 2$ and $E = 1$ eV. It is found that δ_B increases with $\log s_0$. The values of δ_B and δ_P is greater for $a=-2$ than for $a=0$

and 2. For $a=-2$, δ_B decreases from 22% to about 11%, whereas δ_P decreases from 22.5% to 11.2% as $\log s_0$ increases from 8 to 12. For $a=0$, δ_B decreases from 0.3% to 0.17 11% whereas δ_P decreases from 0.6% to 0.34% as $\log s_0$ increases from 8 to 12. For $a=2$,

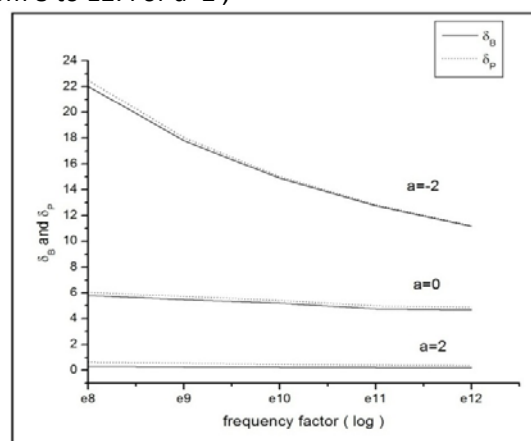


Fig. 3 : Variation of δ_B and δ_P with $\log s_0$ for different values of ‘a’ in Computer simulated glow curve (E=1 eV, b=2)

decreases from 0.6% to 0.34% as $\log s_0$ increases from 8 to 12. For $a=2$, δ_B decreases from 5.8% to about 4.68% whereas δ_P decreases from 6% to 4.86% with the same variation of $\log s_0$. It is to be noted that both Booth method and that due to Pagonis et al [4] underestimate the value of E . We now adjudge the suitability of our findings by considering experimental 376 K TL peak of Ba_{0.92}SO₄:Dy_{0.04}Mn_{0.04} annealed at 873 K and irradiated with 300 Gy γ -ray dose [5]. The values of E_B and E_P are presented in Table 1, along with the values E_{cf} obtained by rigorous curve fitting method[2]. If we consider E_{cf} as the standard or reference value of activation energy, it may be



noted that in accordance with our findings both E_B and E_P are larger than E_{cf} and the value of E_B and E_P are approximately within the range predicted by us.

Table 1. Proportional errors in the determination of E_B and E_P in eV in experimental case [7]

| E_B | E_P | E_{cf} | δ_B (%) | δ_P (%) |
|-------|-------|----------|----------------|----------------|
| 1.085 | 1.091 | 1.08 | 0.463 | 1.02 |

4. CONCLUSION

In the present work, we have analysed the suitability of the two variants of the two heating rates method for the calculation of activation energy namely one proposed by Booth[3] and the other proposed by Pagonis et al[4] both for temperature independent and temperature dependent frequency factors in TL. The percentile proportional error involved are larger for the case of temperature dependent frequency factor than for temperature independent frequency factor. The applicability of the present findings has been tested by considering an experimental TL peak of $Ba_{0.92}SO_4:Dy_{0.04}Mn_{0.04}$. In future we plan to extend the present work to other heating schemes namely hyperbolic and exponential heating schemes.

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